

# Comparative Analysis on Positive Bias Stress-Induced Instability under High $V_{GS}$ /Low $V_{DS}$ and Low $V_{GS}$ /High $V_{DS}$ in Amorphous InGaZnO Thin-Film Transistors

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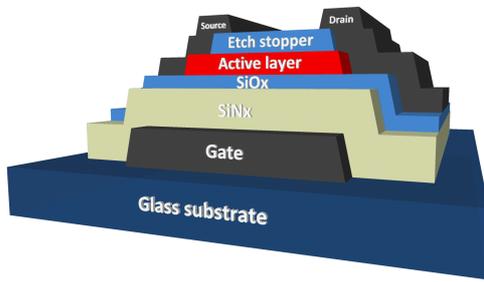
**Abstract**—Positive bias stress-induced instability in amorphous indium-gallium-zinc-oxide (a-IGZO) bottom-gate thin-film transistors (TFTs) was investigated under high  $V_{GS}$ /low  $V_{DS}$  and low  $V_{GS}$ /high  $V_{DS}$  stress conditions through incorporating a forward/reverse  $V_{GS}$  sweep and a low/high  $V_{DS}$  read-out conditions. Our results showed that the electron trapping into the gate insulator dominantly occurs when high  $V_{GS}$ /low  $V_{DS}$  stress is applied. On the other hand, when low  $V_{GS}$ /high  $V_{DS}$  stress is applied, it was found that holes are uniformly trapped into the etch stopper and electrons are locally trapped into the gate insulator simultaneously. During a recovery after the high  $V_{GS}$ /low  $V_{DS}$  stress, the trapped electrons were detrapped from the gate insulator. In the case of recovery after the low  $V_{GS}$ /high  $V_{DS}$  stress, it was observed that the electrons in the gate insulator diffuse to a direction toward the source electrode and the holes were detrapped to out of the etch stopper. Also, we found that the potential profile in the a-IGZO bottom-gate TFT becomes complicatedly modulated during the positive  $V_{GS}/V_{DS}$  stress and the recovery causing various threshold voltages and subthreshold swings under various read-out conditions, and this modulation needs to be fully considered in the design of oxide TFT-based active matrix organic light emitting diode display backplane.

**Index Terms**—a-IGZO TFT, driving condition of AMOLED, positive bias stress, charge trapping, electron-hole pair generation

## I. INTRODUCTION

Amorphous indium-gallium-zinc-oxide (a-IGZO) thin-film transistor (TFT) is actively in use as a promising display device due to many attractive advantages as follows. The a-IGZO TFT has a high mobility along with the good large-area uniformity which is suitable for large-area, high frame-rate, and high resolution display backplane compared with amorphous Si TFT or organic TFT. Moreover, the a-IGZO TFT enables flexible display and transparent display due to its compatibility with a low temperature process and transparency in the visible light range while it still has a high on/off current ratio and good subthreshold swing (SS) [1, 2].

Despite the fact that the a-IGZO TFT enhance displaying industries with its various advantages, the stability under a bias stress still remains as a challenging issue, especially from the viewpoint of mass production. The positive bias stress (PBS)-induced instability is of special importance in terms of the active matrix organic light emitting diode (AMOLED) displays because applying both the positive gate-to-source voltage ( $V_{GS}$ ) and positive drain-to-source voltage ( $V_{DS}$ ) corresponds with the driving condition of AMOLED [3-5]. Although various  $V_{GS}/V_{DS}$  conditions including the high  $V_{DS}$  need to be investigated for optimizing the driving scheme of the AMOLED backplane driven by a-IGZO TFTs, most of the previous works have been focused on the cases of



**Fig. 1.** The a-IGZO TFT with the inverted staggered bottom-gate top-ES structure

$V_{GS}$  higher than or equal to  $V_{DS}$  [6-8].

In this paper, we investigated the PBS instability of the a-IGZO TFT with an inverted staggered bottom-gate top-etch stopper under two stress conditions, i.e., high  $V_{GS}$ /low  $V_{DS}$  and low  $V_{GS}$ /high  $V_{DS}$ , and then compared them each other combining the forward/reverse  $V_{GS}$  sweeps and the high/low  $V_{DS}$  read-out conditions. The main method used in this work was to reproduce the measured electrical characteristics via technology computer-aided-design (TCAD) device simulation. As a consequence, it was verified that the electron trapping into the gate insulator (GI) was dominating in the high  $V_{GS}$ /low  $V_{DS}$  stress whereas the local electron trapping into GI near the drain and hole trapping into etch stopper (ES) layer were dominant in low  $V_{GS}$ /high  $V_{DS}$  stress.

Consequently, we strongly believe that our results can contribute to enhancing and optimizing the driving scheme of a-IGZO TFTs-driven AMOLED from the viewpoint of PBS-induced instability.

## II. DEVICE STRUCTURE, ELECTRICAL CHARACTERISTICS, AND MEASUREMENT SETUP

The a-IGZO TFTs with the inverted staggered bottom-gate top-ES structure were fabricated on a glass substrate as shown in Fig. 1. The detailed process is as follows. Molybdenum electrode was deposited on the glass substrate by using a radio frequency sputtering process and patterned by dry etching. Then, the GI made up of 400 nm/50 nm-thick SiNx/SiOx bilayer was deposited by plasma-enhanced chemical-vapor-deposition (PECVD). The 50-nm-thick a-IGZO thin-film was dc sputter-deposited with power of 3 kW at room temperature in gas mixture of Ar/O<sub>2</sub>=35/63 at sccm in which its atomic ratio

is In:Ga:Zn=2:2:1. The active layer was then patterned by wet etching with diluted HF. Subsequently, the 50-nm-thick SiO<sub>x</sub> was formed as the ES layer by PECVD and dry etch patterning. Molybdenum was then dc sputter-deposited and patterned by dry etching as the source/drain electrodes. Subsequently, the 100-nm-thick SiO<sub>x</sub> was the PECVD-deposited as a passivation layer. Lastly, the post annealing process was performed at 250°C for 1 hour.

The fabricated TFT has geometrical parameters, such as an equivalent oxide thickness of GI  $T_{OX}$ =258 nm, the active layer thickness  $T_{IGZO}$ =50 nm, the channel length  $L$ =50  $\mu$ m, the gate-to-S/D overlap length  $L_{OV}$ =15  $\mu$ m, and the channel width  $W$ =25  $\mu$ m. The electrical properties of the a-IGZO TFT are the threshold voltage ( $V_T$ )=0.4 V,  $SS$ =387 mV/dec, field-effect mobility  $\mu_{FE,lin}$ =7.9 cm<sup>2</sup>/Vs (linear region), and  $\mu_{FE,sat}$ =11.06 cm<sup>2</sup>/Vs (saturation region).

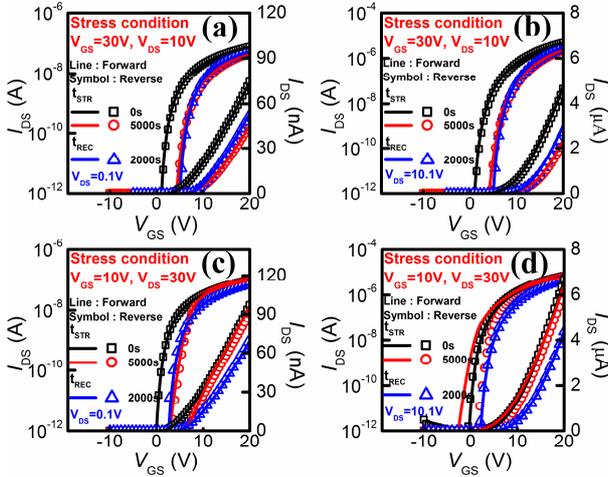
The current-voltage transfer curves were measured at room temperature in dark ambient with the Agilent 4156C precision semiconductor parameter analyzer while the Atlas-2D was used for TCAD device simulation [13].

The PBS was measured under two stress conditions. The high  $V_{GS}$ /low  $V_{DS}$  stress was measured under 30 V/10 V whereas the low  $V_{GS}$ /high  $V_{DS}$  was measured under 10 V/30 V during the same stress time  $t_{STR}$ =5000 sec. In addition, the recovery measurement was conducted during the recovery time, which was set up as  $t_{REC}$ =2000 sec and all electrodes were tied to a ground during the  $t_{REC}$ .

Also, the transfer curve was measured with two kinds of  $V_{GS}$  sweep modes, i.e., forward and reverse  $V_{GS}$  sweeps, and with two kinds of  $V_{DS}$  conditions, i.e., low and high  $V_{DS}$  read-outs, respectively.

The forward  $V_{GS}$  sweep was to set the positions of source and drain electrodes during the read-out of transfer curve to the same as those of during the PBS. Contrastively, the reverse  $V_{GS}$  sweep signifies exchanging the positions of source and drain electrodes during the read-out in comparison with those of during the PBS.

On the other hand, the high  $V_{DS}$ /low  $V_{DS}$  read-out transfer curves were acquired under  $V_{DS}$ =10.1/0.1 V conditions, which individually corresponds to the saturation and linear regions.

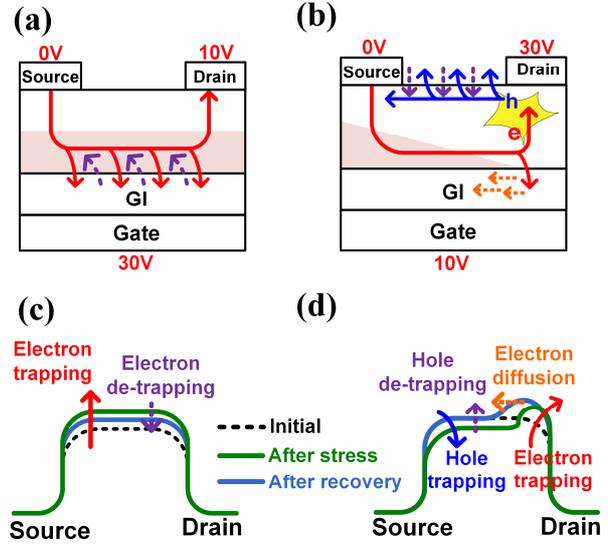


**Fig. 2.** The PBS and recovery time-evolutions of measured transfer characteristics (a) The  $V_{DS}=0.1$  V (linear region) read-out under high  $V_{GS}$ /low  $V_{DS}$  PBS, (b) the  $V_{DS}=10.1$  V (saturation region) read-out under high  $V_{GS}$ /low  $V_{DS}$  PBS, (c) the  $V_{DS}=0.1$  V (linear region) read-out under low  $V_{GS}$ /high  $V_{DS}$  PBS, (d) the  $V_{DS}=10.1$  V (saturation region) read-out under low  $V_{GS}$ /high  $V_{DS}$  PBS

### III. RESULT & DISCUSSION

Fig. 2 shows the measured transfer curves. It is noteworthy that the  $t_{STR}/t_{REC}$ -evolutions of transfer curves were observed in complicated and different ways depending on all of the PBS ( $V_{GS}/V_{DS}=30$  V/10 V and  $V_{GS}/V_{DS}=10$  V/30 V),  $V_{GS}$  sweep (forward and reverse), and  $V_{DS}$  read-out ( $V_{DS}=10.1/0.1$  V) conditions. These time-dependences of transfer curves need to be carefully investigated to understand the instability mechanisms in the operation of IGZO-driven AMOLED.

As shown in Fig. 2(a) and (b), the case of  $V_{GS}/V_{DS}=30$  V/10 V stress condition, the positive  $V_T$  shift is observed during  $t_{STR}$  regardless of the  $V_{DS}$  read-out conditions. Moreover, there is no observed difference between the forward and reverse  $V_{GS}$  sweeps. Based upon such result, it can be inferred that the local  $V_T$  near source is identically increased to that of the drain during the  $t_{STR}$ . These phenomena have been frequently observed by other research groups and analyzed by uniformly electron trapping into the GI or the interface between GI and active layer [10, 11]. Similarly, in our case, electrons can be energetically enhanced by a high vertical electric field and be uniformly trapped into the GI as indicted by solid red arrows in Fig. 3(a), raising the electron potential profile during the  $t_{STR}$  (solid red arrow in Fig. 3(c)).



**Fig. 3.** Schematics illustrating the PBS instability during  $t_{STR}$  and  $t_{REC}$  under (a) high  $V_{GS}$ /low  $V_{DS}$ , (b) low  $V_{GS}$ /high  $V_{DS}$  PBS conditions. Schematics illustrating the lateral energy band diagrams and electron potential profiles during  $t_{STR}$  and  $t_{REC}$  under (a) high  $V_{GS}$ /low  $V_{DS}$ , (b) low  $V_{GS}$ /high  $V_{DS}$  PBS conditions

Consequently, the positive  $V_T$  shift without SS degradation is observed in all cases of the read-out conditions. The transfer curve is also slightly shifted to a negative  $V_{GS}$  direction during the  $t_{REC}$  as shown in Fig. 2(a) and (b), which can be attributed to the electron detrapping from GI as illustrated by dashed purple arrows in Fig. 3(a) and (c).

However, in term of  $V_{GS}/V_{DS}=10$  V/30 V stress condition, the measured transfer curves obviously have different  $t_{STR}/t_{REC}$ -dependences in each read-out condition as represented in Fig. 2(c) and (d). In the case of the  $V_{DS}=0.1$  V read-out condition, the difference between the forward and reverse  $V_{GS}$  sweeps was negligible and both the SS degradation and the slightly positive  $V_T$  shift with hump phenomenon were clearly observed after  $t_{STR}=5000$  sec as shown in Fig. 2(c). On the contrary, in case of the  $V_{DS}=10.1$  V read-out condition, there is definite distinction between the forward and reverse  $V_{GS}$  sweeps. The positive  $V_T$  shift with SS degradation is shown in reverse sweep; however, the forward  $V_{GS}$  sweep has the result of negative  $V_T$  shift with SS degradation as shown in Fig. 2(d). That is to say, opposite direction of the  $V_T$  shift is caused only under the condition of  $V_{DS}=10.1$  V read-out (Fig. 2(d)), and this phenomenon indicates the asymmetric  $V_T$  is formed in

each near the source and drain after the stress. Additionally, SS degradation is the clue to predict occurring of the asymmetric degradation [9]. We have come to understand abovementioned phenomena as follows. The local  $V_T$  near the source is lowered whereas the local  $V_T$  near the drain becomes higher during  $t_{STR}$  as a consequence of the local electron trapping into GI near the drain. In the low  $V_{GS}$ /high  $V_{DS}$  stress condition, the high local lateral electric field is formed near the drain because of high drain voltage. Therefore, electron, the mobile carrier, is rapidly accelerated by the high electric field and then can obtain the large kinetic energy, which is large enough to generate the electron-hole pair caused by the impact ionization [7, 12] or by the interaction between the valence band and subgap traps under mid-gap of a-IGZO. Accordingly, the electrons can be locally trapped into GI only near the drain (solid red arrow in Fig. 3(b)); hence, the local  $V_T$  as well as the electron potential near the drain becomes higher than the initial (solid red arrow in Fig. 3(d)). Generated holes can drift to the source under the influence of the lateral electric field so they can be quite uniformly trapped into ES (solid blue arrow in Fig. 3(b)), i.e., the  $V_T$  near the source becomes lower during the  $t_{STR}$  (solid blue arrow in Fig. 3(d)). Consequently, the lateral energy band of active layer is asymmetrically formed as indicated in Fig. 3(d) after the low  $V_{GS}$ /high  $V_{DS}$  stress. This phenomenon could be distinguished only in  $V_{DS}=10.1$  V read-out and it can bring about the different characteristic in each read-out condition because each condition has different  $V_T$  depending not only on the forward and reverse  $V_{GS}$  sweeps but also how high  $V_{DS}$  is applied. When correlating between the before-mentioned analysis and measured data, the negative  $V_T$  shift is observed only under the forward high  $V_{DS}$  read-out because this condition can completely remove the locally high electron potential barrier near the drain due to applying high  $V_{DS}$ .

After  $t_{REC}=2000$  s under  $V_{GS}/V_{DS}=10$  V/30 V stress condition, we could see all read-out transfer curves move towards positive  $V_{GS}$  direction. Especially, in the case of measured transfer curve under the forward and  $V_{DS}=10.1$  V read-out condition, the negative  $V_T$  shift after the stress showed the complete recovery. More specifically, the transfer curve moved more towards the positive direction compared to the initial condition, that

is to say, there is no difference between the forward and reverse  $V_{GS}$  sweeps. Therefore, it should be deduced that the recovery characteristic is caused by hole detrapping (dashed purple arrows in Fig. 3(b) and (d)). Furthermore, we consider the disappearance of the difference between the forward and reverse  $V_{GS}$  sweep transfer curves were consequent on the electric field-enhanced diffusion of trapped electrons in GI towards the source due not only to the electron concentration difference in GI, but also to the built-in electric field resulting from the asymmetric potential distribution as well (dashed orange arrows in Fig. 3(b) and (d)). However, the electron movement during the  $t_{REC}$  needs to be further investigated and proven.

#### IV. TCAD SIMULATION

In Fig. 3, we established the qualitative model on the  $t_{STR}/t_{REC}$ -dependence of transfer curves in all cases of PBS,  $V_{GS}$  sweep, and  $V_{DS}$  read-out conditions. In addition, we conducted the TCAD device simulation in order to validate our model quantitatively. For a precise and reasonable analysis, the subgap trap model of a-IGZO TFT [14] was employed as follows:

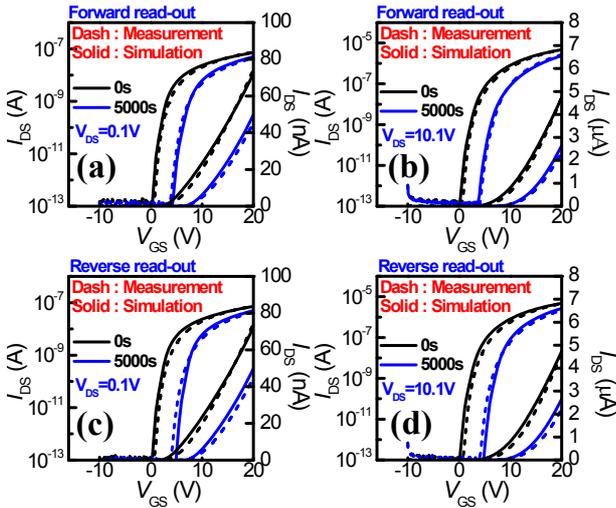
$$\begin{aligned} g_A(E) &= g_{DA}(E) + g_{TA}(E) \\ &= N_{DA} \times \exp\left(\frac{E - E_C}{kT_{DA}}\right) + N_{TA} \times \exp\left(\frac{E - E_C}{kT_{TA}}\right) \\ g_{TD}(E) &= N_{TD} \times \exp\left(\frac{E_V - E}{kT_{TD}}\right). \end{aligned}$$

$g_A(E)$ , which indicates the acceptor-like states near the conduction band, is composed of exponential tail states ( $N_{DA}$ ,  $kT_{DA}$ ) and exponential deep states ( $N_{TA}$ ,  $kT_{TA}$ ). Also,  $g_D(E)$  indicating the donor-like states near the valence band is expressed by an exponential function ( $N_{TD}$ ,  $kT_{TD}$ ). The parameters utilized in the TCAD simulation are summarized in Table 1 [15]. The TCAD parameters used in simulation are reasonable because the simulated transfer curves accurately correspond to the measured initial transfer curves (black solid and black dashed lines in Figs. 4 and 5).

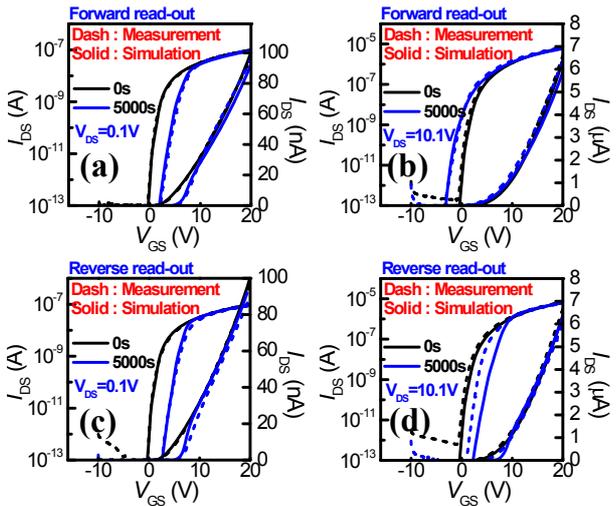
In order to verify our proposed model in  $V_{GS}/V_{DS}=30$  V/10 V stress condition, simulation was conducted by uniformly putting the electron fixed charge in GI. As a

**Table 1.** The parameters used in TCAD simulation

$E_g$ [eV]	$N_C$ [cm <sup>-3</sup> ]	$N_V$ [cm <sup>-3</sup> ]	$N_{SD}$ [cm <sup>-3</sup> ]	$\mu_n$ [cm <sup>2</sup> /Vs]	
3.2	1x10 <sup>18</sup>	2x10 <sup>19</sup>	1.2x10 <sup>16</sup>	55	
$N_{TA}$ [cm <sup>-3</sup> eV <sup>-1</sup> ]	$kT_{TA}$ [eV]	$N_{DA}$ [cm <sup>-3</sup> eV <sup>-1</sup> ]	$kT_{DA}$ [eV]	$N_{TD}$ [cm <sup>-3</sup> eV <sup>-1</sup> ]	$kT_{TD}$ [eV]
9x10 <sup>19</sup>	0.02	1x10 <sup>18</sup>	0.16	1x10 <sup>18</sup>	0.27



**Fig. 4.** Measured and simulated transfer curves under high  $V_{GS}/low V_{DS}$  PBS (a) The forward  $V_{DS}=0.1$  V read-out, (b) the forward  $V_{DS}=10.1$  V read-out, (c) the reverse  $V_{DS}=0.1$  V read-out, (d) the reverse  $V_{DS}=10.1$  V read-out



**Fig. 5.** Measured and simulated transfer curves under low  $V_{GS}/high V_{DS}$  PBS (a) The forward  $V_{DS}=0.1$  V read-out, (b) the forward  $V_{DS}=10.1$  V read-out, (c) the reverse  $V_{DS}=0.1$  V read-out, (d) the reverse  $V_{DS}=10.1$  V read-out

result, simulated transfer curves were well reproduced under all case of read-out conditions that is identical to measurement transfer curves such as solid blue lines and dashed blue lines in Fig. 4. Moreover, we were able to

figure out that trapped electron charge density during the  $t_{STR}=5000$  sec is  $4 \times 10^{11}$  cm<sup>-2</sup> which is in the range of reasonable value.

In order to verify our model in  $V_{GS}/V_{DS}=10$  V/30 V stress condition, the electrons were input from the drain edge to the source direction with a distance of 1  $\mu$ m in the GI and hole were uniformly put into ES from the source edge to drain edge. Consequently, it should be emphasized that there is verified correlation between the measured data and simulated data. Also, we find out that  $4.7 \times 10^{11}$  cm<sup>-2</sup> of electrons and  $2 \times 10^{11}$  cm<sup>-2</sup> of holes were trapped respectively during the  $t_{STR}=5000$  sec.

Through the TCAD simulation, measurement result is reproduced under the two stress conditions which are high  $V_{GS}/low V_{DS}$  and low  $V_{GS}/high V_{DS}$  with all measurement read-out conditions. Therefore, it is proved that our instability model is reasonable.

### V. CONCLUSIONS

We have analyzed the PBS-induced instability under high  $V_{GS}/low V_{DS}$  and low  $V_{GS}/high V_{DS}$  conditions in the bottom gate a-IGZO TFTs with an emphasis on the forward/reverse and a low/high  $V_{DS}$  read-out conditions. In case of high  $V_{GS}/low V_{DS}$  stress condition, the electrons are uniformly trapped into gate insulator by the vertical electric field. On the other hands, in case of low  $V_{GS}/high V_{DS}$ , it was dominated by the local trapping of electron/hole into GI/ES and the local potential barriers were complicatedly modulated by the variations of not only the position of the source/drain but also the value of  $V_{DS}$  under an individual read-out condition. The proposed instability model was meticulously verified by TCAD simulation under all the stress conditions and read-outs.

During the  $t_{REC}$ , it was observed and analyzed that electron detrapping occurs after high  $V_{GS}/low V_{DS}$  stress whereas electron diffusion and hole detrapping primarily occurs after low  $V_{GS}/high V_{DS}$  stress; however, further studies should be required.

Consequently, we believe that our results would provide the informative implications of optimizing the driving scheme for the display pixel in the amorphous oxide TFT-driven AMOLED backplanes.

## ACKNOWLEDGMENTS

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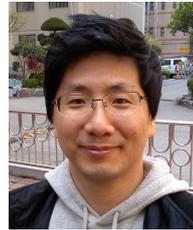
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